

### 11.1.2 Electric Dipole Radiation

$$R_{\pm} = \sqrt{r^2 \mp rd \cos \theta + (d/2)^2} = r \left[ 1 \mp \frac{d}{r} \cos \theta + \frac{1}{r^2} \left( \frac{d}{2} \right)^2 \right]^{\frac{1}{2}} \quad (11.6)$$

approx.1  $d \ll r$

$$R_{\pm} \cong r \left( 1 \mp \frac{d}{2r} \cos \theta \right)$$

approx.2  $d \ll \frac{c}{\omega} = \frac{\lambda}{2\pi} = \frac{1}{k}$

$$\cos[\omega(t - R_{\pm}/c)] \cong \cos[\omega(t - r/c)] \cos\left(\pi \frac{d}{\lambda} \cos \theta\right) \mp \sin[\omega(t - r/c)] \sin\left(\pi \frac{d}{\lambda} \cos \theta\right)$$

$$\cos[\omega(t - R_{\pm}/c)] \cong \cos[\omega(t - r/c)] \mp \pi \frac{d}{\lambda} \cos \theta \sin[\omega(t - r/c)] \quad (11.11)$$

(11.5)

$$\begin{aligned} V(\mathbf{r}, t) &= \frac{1}{4\pi\epsilon_0} \frac{q_0}{r} \left\{ \left(1 + \frac{d}{2r} \cos \theta\right) \cos[\omega(t - R_+/c)] - \left(1 - \frac{d}{2r} \cos \theta\right) \cos[\omega(t - R_-/c)] \right\} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q_0}{r} \left\{ \left(1 + \frac{d}{2r} \cos \theta\right) \left[ \cos[\omega(t - r/c)] - \pi \frac{d}{\lambda} \cos \theta \sin[\omega(t - r/c)] \right] \right. \\ &\quad \left. - \left(1 - \frac{d}{2r} \cos \theta\right) \left[ \cos[\omega(t - r/c)] + \pi \frac{d}{\lambda} \cos \theta \sin[\omega(t - r/c)] \right] \right\} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q_0}{r} \left\{ \frac{d}{r} \cos \theta \cos[\omega(t - r/c)] - 2\pi \frac{d}{\lambda} \cos \theta \sin[\omega(t - r/c)] \right\} \\ &= \frac{q_0 d \cos \theta}{4\pi\epsilon_0 r} \left\{ \frac{1}{r} \cos[\omega(t - r/c)] - \frac{2\pi}{\lambda} \sin[\omega(t - r/c)] \right\} \quad (11.12) \end{aligned}$$

approx.3  $r \gg \frac{c}{\omega} = \frac{\lambda}{2\pi} = \frac{1}{k}$

$$V(r, \theta, t) = \frac{q_0 d \cos \theta}{4\pi\epsilon_0 r} \left\{ -\frac{2\pi}{\lambda} \sin[\omega(t - r/c)] \right\} \quad (11.14)$$

$$A(r, t) = \frac{\mu_0}{4\pi} \int_{-d/2}^{d/2} \frac{-q_0 \omega \sin[\omega(t - R/c)] \hat{z}}{R} dz \quad (11.16) \quad \text{被積分関数は} z \text{に依存}$$

$$= \hat{z} \frac{\mu_0}{4\pi} (-q_0 \omega) \int_{-d/2}^{d/2} \frac{1}{r} \left(1 + \frac{d}{2r} \cos \theta\right) \sin\left[\omega t - \frac{\omega}{c} r \left(1 - \frac{d}{2r} \cos \theta\right)\right] dz$$

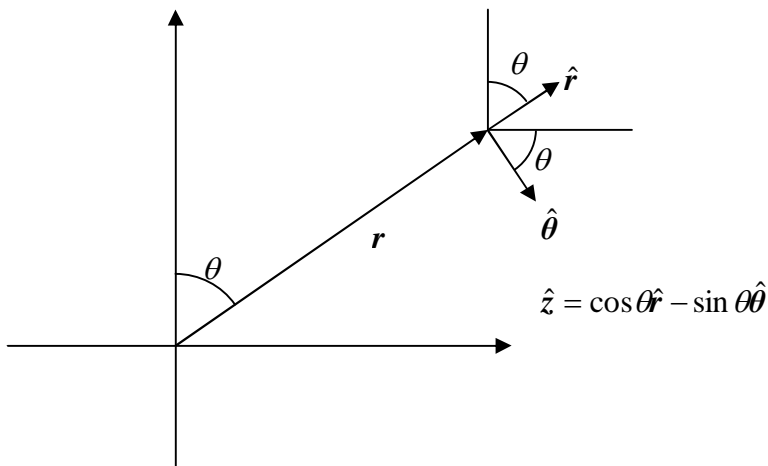
$$= \hat{z} \frac{\mu_0}{4\pi} (-q_0 \omega) \int_{-d/2}^{d/2} \frac{1}{r} \left(1 + \frac{d}{2r} \cos \theta\right) \sin\left[\omega\left(t - \frac{r}{c}\right) + \pi \frac{d}{\lambda} \cos \theta\right] dz$$

approx.1  $d \ll r$       approx.2  $d \ll \frac{c}{\omega} = \frac{\lambda}{2\pi} = \frac{1}{k}$

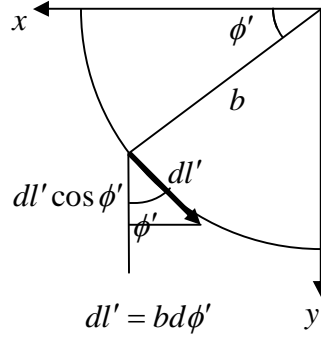
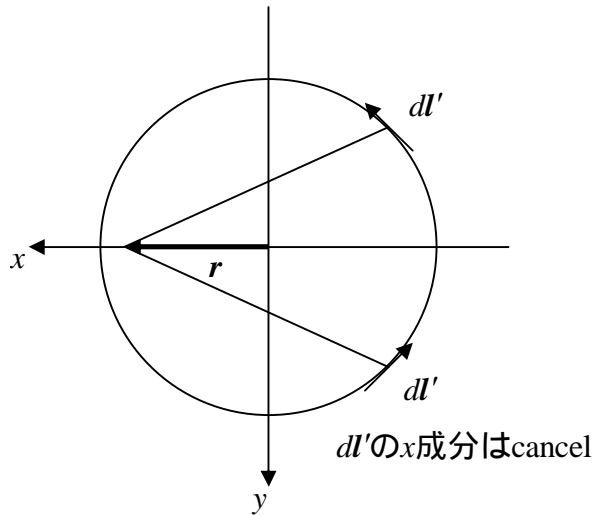
積分によって  $d$  の1次のオーダーになるので、被積分関数は  $d$  について1次の項は落とす

$$= \hat{z} \frac{\mu_0}{4\pi} (-q_0 \omega) \int_{-d/2}^{d/2} \frac{1}{r} \sin\left[\omega\left(t - \frac{r}{c}\right)\right] dz \quad \text{被積分関数は} z \text{に依存しない}$$

$$= -\frac{\mu_0 q_0 d \omega}{4\pi r} \sin\left[\omega\left(t - \frac{r}{c}\right)\right] \hat{z} \quad (11.17)$$



### 11.1.3 Magnetic Dipole Radiation



(11.27)に(11.30)(11.32)を代入

$$\begin{aligned} \mathbf{A}(\mathbf{r}, t) &= \frac{\mu_0 I_0 b}{4\pi} \hat{\mathbf{y}} \int_0^{2\pi} \frac{\cos[\omega(t - R/c)]}{R} \cos \phi' d\phi' \\ &= \frac{\mu_0 I_0 b}{4\pi} \hat{\mathbf{y}} \int_0^{2\pi} \frac{1}{r} \left(1 + \frac{b}{r} \sin \theta \cos \phi'\right) \left\{ \cos[\omega(t - r/c)] - \frac{\omega b}{c} \sin \theta \cos \phi' \sin[\omega(t - r/c)] \right\} \cos \phi' d\phi' \end{aligned}$$

2nd - order term  $\frac{b}{r} \cdot \frac{\omega b}{c}$  の項を無視

(11.35)  $\rightarrow$  (11.37)

$$\mathbf{A}(r, \theta, t) = A_\phi(r, \theta, t) \hat{\boldsymbol{\phi}}$$

$$\nabla \times \mathbf{A} = \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (A_\phi \sin \theta) \hat{\mathbf{r}} + \frac{1}{r} \left\{ -\frac{\partial}{\partial r} (r A_\phi) \right\} \hat{\boldsymbol{\theta}}$$

$$\text{第1項} \propto \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} \left\{ \frac{\sin^2 \theta}{r} \sin[\omega(t - r/c)] \right\} = \frac{1}{r \sin \theta} \frac{2 \sin \theta \cos \theta}{r} \sin[\omega(t - r/c)] = \frac{1}{r} \frac{2 \cos \theta}{r} \sin[\omega(t - r/c)]$$

$$\text{第2項} \propto \frac{1}{r} \left\{ -\frac{\partial}{\partial r} (\sin \theta \sin[\omega(t - r/c)]) \right\} = \frac{1}{r} \frac{\omega}{c} \sin \theta \cos[\omega(t - r/c)]$$

approx.3  $\frac{1}{r} \ll \frac{\omega}{c}$  第1項を無視 第2項が(11.37)

### 11.1.4 Radiation from an Arbitrary Source

$$(11.50) \quad \dot{\rho}\left(\frac{\hat{\mathbf{r}} \cdot \mathbf{r}'}{c}\right) \gg \frac{1}{2} \ddot{\rho}\left(\frac{\hat{\mathbf{r}} \cdot \mathbf{r}'}{c}\right)^2 \quad \dot{\rho} \gg \ddot{\rho}\left(\frac{\hat{\mathbf{r}} \cdot \mathbf{r}'}{c}\right) \quad \frac{c}{\dot{\rho} / \ddot{\rho}} \gg \hat{\mathbf{r}} \cdot \mathbf{r}'$$

$$V(\mathbf{r}, t) \cong \frac{1}{4\pi\epsilon_0} \int \frac{1}{r} \left(1 + \frac{\hat{\mathbf{r}} \cdot \mathbf{r}'}{r}\right) [\rho(\mathbf{r}', t_0) + \dot{\rho}(\mathbf{r}', t_0) \left(\frac{\hat{\mathbf{r}} \cdot \mathbf{r}'}{c}\right)] d\tau' \quad r' \text{ について1次まで残す}$$

p.214 Prob.5.7  $\mathbf{p} = \int \mathbf{r}' \rho(\mathbf{r}', t) d\tau' \quad \nabla \cdot (\mathbf{x}\mathbf{J}) = (\nabla \mathbf{x}) \cdot \mathbf{J} + \mathbf{x} \nabla \cdot \mathbf{J} = J_x + \mathbf{x} \left(-\frac{\partial \rho}{\partial t}\right)$

$$\dot{p}_x = \int \mathbf{x}' \frac{\partial \rho}{\partial t} d\tau' = \int J_x d\tau' - \int \nabla \cdot (\mathbf{x}'\mathbf{J}) d\tau' = \int J_x d\tau' - \int (\mathbf{x}'\mathbf{J}) \cdot d\mathbf{S}' = \int J_x d\tau' \quad \therefore \dot{\mathbf{p}} = \int \mathbf{J} d\tau'$$

$$\begin{aligned} \nabla \left[ \frac{\hat{\mathbf{r}} \cdot \dot{\mathbf{p}}(t_0)}{r} \right] &= \nabla \left[ \frac{\mathbf{r} \cdot \dot{\mathbf{p}}(t_0)}{r^2} \right] = \nabla \left[ \frac{x\dot{p}_x + y\dot{p}_y + z\dot{p}_z}{x^2 + y^2 + z^2} \right] \\ &= \left( \mathbf{i} \frac{\partial}{\partial x} + \mathbf{j} \frac{\partial}{\partial y} + \mathbf{k} \frac{\partial}{\partial z} \right) \left[ \frac{x\dot{p}_x + y\dot{p}_y + z\dot{p}_z}{x^2 + y^2 + z^2} \right] \\ &= \mathbf{i} \left[ \frac{\partial}{\partial x} \left( \frac{1}{x^2 + y^2 + z^2} \right) (x\dot{p}_x + y\dot{p}_y + z\dot{p}_z) + \left( \frac{1}{x^2 + y^2 + z^2} \right) \frac{\partial}{\partial x} (x\dot{p}_x + y\dot{p}_y + z\dot{p}_z) \right] + \mathbf{j}[\ ] + \mathbf{k}[\ ] \end{aligned}$$

$$\frac{\partial}{\partial x} \left( \frac{1}{x^2 + y^2 + z^2} \right) = -\frac{2x}{(x^2 + y^2 + z^2)^2} = -\frac{2x}{r^4} \quad \text{この項は無視できる}$$

残るのは

$$\frac{1}{x^2 + y^2 + z^2} \left[ \mathbf{i} \frac{\partial}{\partial x} (x\dot{p}_x + y\dot{p}_y + z\dot{p}_z) + \mathbf{j} \frac{\partial}{\partial y} (x\dot{p}_x + y\dot{p}_y + z\dot{p}_z) + \mathbf{k} \frac{\partial}{\partial z} (x\dot{p}_x + y\dot{p}_y + z\dot{p}_z) \right]$$

$$\begin{aligned} \nabla \left[ \frac{\hat{\mathbf{r}} \cdot \dot{\mathbf{p}}(t_0)}{r} \right] &= \nabla \left[ \frac{\mathbf{r} \cdot \dot{\mathbf{p}}(t_0)}{r^2} \right] \cong \frac{1}{r^2} \nabla (\mathbf{r} \cdot \dot{\mathbf{p}}(t_0)) \\ &= \frac{1}{r^2} \left[ \mathbf{i} \left( \dot{p}_x + x \frac{\partial}{\partial x} \dot{p}_x + y \frac{\partial}{\partial x} \dot{p}_y + z \frac{\partial}{\partial x} \dot{p}_z \right) + \mathbf{j} \left( \dot{p}_y + x \frac{\partial}{\partial y} \dot{p}_x + y \frac{\partial}{\partial y} \dot{p}_y + z \frac{\partial}{\partial y} \dot{p}_z \right) + \mathbf{k} \left( \dot{p}_z + x \frac{\partial}{\partial z} \dot{p}_x + y \frac{\partial}{\partial z} \dot{p}_y + z \frac{\partial}{\partial z} \dot{p}_z \right) \right] \end{aligned}$$

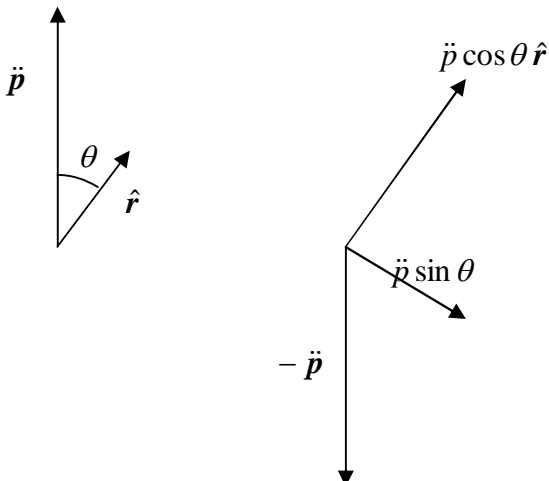
$\frac{\dot{p}_x}{r^2}, \frac{\dot{p}_y}{r^2}, \frac{\dot{p}_z}{r^2}$  を無視

$$\cong \frac{1}{r^2} \left[ \mathbf{i} \left( x \frac{\partial t_0}{\partial x} \frac{\partial \dot{p}_x}{\partial t_0} + y \frac{\partial t_0}{\partial x} \frac{\partial \dot{p}_y}{\partial t_0} + z \frac{\partial t_0}{\partial x} \frac{\partial \dot{p}_z}{\partial t_0} \right) + \mathbf{j}(\ ) + \mathbf{k}(\ ) \right]$$

$$= \frac{1}{r^2} \left[ \mathbf{i} \left( x \frac{\partial \dot{p}_x}{\partial t_0} + y \frac{\partial \dot{p}_y}{\partial t_0} + z \frac{\partial \dot{p}_z}{\partial t_0} \right) \frac{\partial t_0}{\partial x} + \mathbf{j}(\ ) + \mathbf{k}(\ ) \right]$$

$$= \frac{1}{r^2} (\mathbf{r} \cdot \ddot{\mathbf{p}}) \nabla t_0$$

$$\nabla \times \dot{\mathbf{p}}(t_0) \quad x \text{成分} \quad \frac{\partial \dot{p}_z}{\partial y} - \frac{\partial \dot{p}_y}{\partial z} = \frac{\partial t_0}{\partial y} \frac{\partial \dot{p}_z}{\partial t_0} - \frac{\partial t_0}{\partial z} \frac{\partial \dot{p}_y}{\partial t_0}$$



$$\begin{aligned}
\nabla \frac{Q(t-r/c)}{r} &= \frac{1}{r} \nabla Q + Q \nabla \frac{1}{r} \\
&= \frac{1}{r} \frac{\partial t_0}{\partial r} \frac{dQ}{dt_0} - \frac{Q}{r^2} \hat{\mathbf{r}} \\
&= -\frac{1}{cr} \frac{\mathbf{r}}{r} \dot{Q} - \frac{Q}{r^2} \hat{\mathbf{r}} \\
&= -\frac{\dot{Q}}{cr} \hat{\mathbf{r}} - \frac{Q}{r^2} \hat{\mathbf{r}}
\end{aligned}$$

## 11.2 Point Charges

(11.66) → (11.68)

$$\mathbf{E}_{\text{rad}} = \frac{q}{4\pi\epsilon_0} \frac{R}{(\mathbf{R} \cdot \mathbf{u})^3} [\mathbf{R} \times (\mathbf{u} \times \mathbf{a})] = \frac{q}{4\pi\epsilon_0} \frac{R}{(cR)^3} [\mathbf{R} \times (c\hat{\mathbf{R}} \times \mathbf{a})] = \frac{q}{4\pi\epsilon_0 c^2 R} [\hat{\mathbf{R}} \times (\hat{\mathbf{R}} \times \mathbf{a})] = \frac{q}{4\pi\epsilon_0 c^2 R} [(\hat{\mathbf{R}} \cdot \mathbf{a})\hat{\mathbf{R}} - (\hat{\mathbf{R}} \cdot \hat{\mathbf{R}})\mathbf{a}]$$

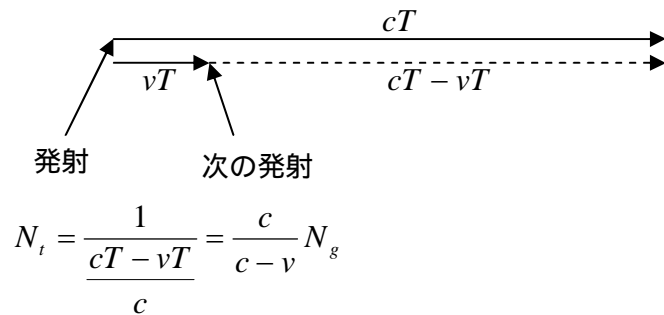
$$\because \mathbf{u} = c\hat{\mathbf{R}} \quad \mathbf{R} \cdot \mathbf{u} = cR$$

(11.68) → (11.69)

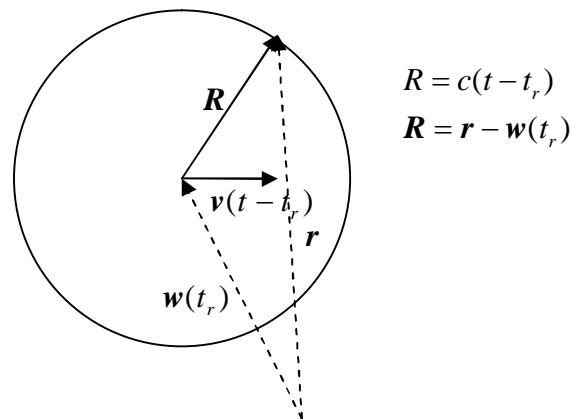
$$[(\hat{\mathbf{R}} \cdot \mathbf{a})\hat{\mathbf{R}} - \mathbf{a}] \cdot [(\hat{\mathbf{R}} \cdot \mathbf{a})\hat{\mathbf{R}} - \mathbf{a}] = (\hat{\mathbf{R}} \cdot \mathbf{a})^2 + a^2 - 2(\hat{\mathbf{R}} \cdot \mathbf{a})^2 = a^2 - (\hat{\mathbf{R}} \cdot \mathbf{a})^2$$

$$(11.70) \quad \int_0^\pi \sin^3 \theta d\theta = \frac{4}{3}$$

$$T: \text{発射周期} \quad N_g = \frac{1}{T}$$



$$N_g = \frac{c - v}{c} N_t = \left(1 - \frac{v}{c}\right) N_t = \left(1 - \frac{\hat{\mathbf{R}} \cdot \mathbf{v}}{c}\right) N_t$$



$$(11.66) \quad \mathbf{E}_{\text{rad}} = \frac{q}{4\pi\epsilon_0} \frac{R}{(\mathbf{R} \cdot \mathbf{u})^3} [\mathbf{R} \times (\mathbf{u} \times \mathbf{a})]$$

$$\begin{aligned}
(11.72) \quad \frac{dP}{d\Omega} &= \left(\frac{\mathbf{R} \cdot \mathbf{u}}{Rc}\right) \frac{1}{\mu_0 c} E_{\text{rad}}^2 R^2 = \left(\frac{\mathbf{R} \cdot \mathbf{u}}{Rc}\right) \frac{1}{\mu_0 c} \left\{ \frac{q}{4\pi\epsilon_0} \frac{R}{(\mathbf{R} \cdot \mathbf{u})^3} [\mathbf{R} \times (\mathbf{u} \times \mathbf{a})] \right\}^2 R^2 \\
&= \left(\frac{\mathbf{R} \cdot \mathbf{u}}{Rc}\right) \frac{1}{\mu_0 c} \left(\frac{q}{4\pi\epsilon_0}\right)^2 \frac{R^4}{(\mathbf{R} \cdot \mathbf{u})^6} |\mathbf{R} \times (\mathbf{u} \times \mathbf{a})|^2 \\
&= \frac{1}{\mu_0 c^2} \left(\frac{q}{4\pi\epsilon_0}\right)^2 \frac{R^3}{(\mathbf{R} \cdot \mathbf{u})^5} |\mathbf{R} \times (\mathbf{u} \times \mathbf{a})|^2 \\
&= \epsilon_0 \left(\frac{q^2}{16\pi^2 \epsilon_0^2}\right) \frac{1}{(\hat{\mathbf{R}} \cdot \mathbf{u})^5} |\hat{\mathbf{R}} \times (\mathbf{u} \times \mathbf{a})|^2 = \frac{q^2}{16\pi^2 \epsilon_0} \frac{|\hat{\mathbf{R}} \times (\mathbf{u} \times \mathbf{a})|^2}{(\hat{\mathbf{R}} \cdot \mathbf{u})^5}
\end{aligned}$$